### 2.4 The Entropy

An important characteristic feature for a random variable is the entropy, which we will introduce in this section.

## 2.4.1 Entropy for a Discrete Set of Events

Let  $\{A_1, \ldots, A_N\}$  be a complete, disjoint set of events, i.e.,

$$A_1 \cup A_2 \cup \ldots \cup A_N = \Omega. \tag{2.78}$$

Furthermore, let  $\mathcal P$  be a probability defined for these events. We then define the entropy as

$$S = -k \sum_{i=1} \mathcal{P}(A_i) \ln(\mathcal{P}(A_i)). \tag{2.79}$$

Here k represents a factor which we set equal to 1 for the moment. In the framework of statistical mechanics k will be Boltzmann's constant  $k_{\rm B}$ .

We observe:

- The entropy is defined for a complete, disjoint set of events of a random variable, irrespective of whether this partition of  $\Omega$  into events can be refined or not. If  $\Omega$  is the real axis, we might have, e.g.,  $N=2, A_1=(-\infty,0), A_2=[0,\infty)$ .
- Since  $0 \le \mathcal{P}(A_i) \le 1$  we always have  $S \ge 0$ .
- If  $\mathcal{P}(A_j) = 1$  for a certain j and  $\mathcal{P}(A_i) = 0$  otherwise, then S = 0. This means that if the event  $A_j$  occurs with certainty the entropy is zero.
- If an event has occurred, then, as we will show in a moment,  $-\log_2 \mathcal{P}(A_j)$  is a good measure of the number of questions to be asked in order to find out that it is just  $A_j$  which is realized. In this context, 'question' refers to questions which can be answered by 'yes' or 'no', i.e., the answer leads to a gain of information of 1 bit. Hence, on average the required number of yes-or-no questions is

$$S' = -\sum_{j=1}^{\infty} \mathcal{P}(A_j) \log_2(\mathcal{P}(A_j)) = S + \text{const}.$$
 (2.80)

The entropy is thus a measure of the missing information needed to find out which result is realized.

To show that  $-\log_2 \mathcal{P}(A_j)$  is just equal to the number of required yesor-no questions, we first divide  $\Omega$  into two disjoint domains  $\Omega_1$  and  $\Omega_2$  such that

$$\sum_{A_i \in \Omega_1} \mathcal{P}(A_i) = \sum_{A_i \in \Omega_2} \mathcal{P}(A_i) = \frac{1}{2}.$$
 (2.81)

The first question is now: Is  $A_j$  in  $\Omega_1$ ? Having the answer to this question we next consider the set containing  $A_j$  and multiply the probabilities for the events in this set by a factor of 2. The sum of the probabilities for this set is now again equal to 1, and we are in the same position as before with the set  $\Omega$ : We divide it again and ask the corresponding yes-or-no question. This procedure ends after k steps, where k is the smallest integer such that  $2^k \mathcal{P}(A_j)$  becomes equal to or larger than 1. Consequently,  $-\log_2 \mathcal{P}(A_j)$  is a good measure of the number of yes-or-no questions needed.

• If the probabilities of the events are equal, i.e.,

$$\mathcal{P}(A_i) = \frac{1}{N} \,, \tag{2.82}$$

we have

$$S = \ln N \tag{2.83}$$

Any other distribution of probabilities leads to a smaller S. This will be shown soon.

The above observations suggest that the entropy may be considered as a lack of information when a probability density is given. On average it would require the answers to S yes-or-no questions to figure out which event has occurred. This lack is zero for a density which describes the situation where one event occurs with certainty. If all events are equally probable, this lack of information about which event will occur in a realization is maximal.

A less subjective interpretation of entropy arises when we think of it as a measure for uncertainty. If the probability is the same for all events, the uncertainty is maximal.

# 2.4.2 Entropy for a Continuous Space of Events

In a similar manner we define the entropy for a random variable X, where the space of events is a continuum, by

$$S[\varrho_X] = -k \int dx \, \varrho_X(x) \ln \left( \frac{\varrho_X(x)}{\varrho_0} \right) \,. \tag{2.84}$$

When  $\varrho_X(x)$  has a physical dimension, the denominator  $\varrho_0$  in the argument of the logarithm cannot simply be set to 1. Since the physical dimension of  $\varrho_X(x)$  is equal to the dimension of 1/dx, the physical dimension of  $\varrho_0$  has to be the same, in order that the argument of the logarithm will be dimensionless.

It is easy to see that a change of  $\varrho_0$  by a factor  $\alpha$  leads to a change of the entropy by an additive term  $k \ln \alpha$ . Such a change of  $\varrho_0$  only shifts the scale of S. Notice that we no longer have  $S \geq 0$ .

We now calcula We obtain (for  $k_{
m B}$  -

$$S = \int dx \left( \frac{1}{2} \right)$$
$$= \frac{1}{2} (1 + \ln x)$$

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### 2.4.3 Relative E

The relative entro density function q(

$$S[p|q] = -k$$

Obviously,  $p(x) \equiv$  for a complete and entropy of a densi q(x) is negative se

$$S[p|q] \leq 0.$$

This is easy to see

$$\ln z \le z - 1$$

for 
$$z = \frac{q(x)}{p(x)}$$
, mult

Since both densitiequal, from which

 $\int \mathrm{d}x\,p(x)$ 

#### 2.4.4 Remarks

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